

# HIGH-ORDER UNIFORMLY ACCURATE TIME INTEGRATORS FOR SEMILINEAR WAVE EQUATIONS OF KLEIN–GORDON TYPE IN THE NON-RELATIVISTIC LIMIT

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ABSTRACT. We introduce a family of high-order time semi-discretizations for semilinear wave equations of Klein–Gordon type with arbitrary smooth nonlinearities that are uniformly accurate in the non-relativistic limit where the speed of light goes to infinity. Our schemes do not require pre-computations that are specific to the nonlinearity and have no restrictions in step size. Instead, they rely upon a general oscillatory quadrature rule developed in a previous paper (Mohamad and Oliver, arXiv:1909.04616).

## 1. INTRODUCTION

Let  $X$  be a Hilbert space with inner product  $\langle \cdot, \cdot \rangle$  and associated norm  $\|\cdot\|$ . We study a semilinear wave equation on  $X$ ,

$$c^{-2} \partial_{tt}\phi + L\phi + c^2 \phi = f(\phi, t), \quad (1a)$$

$$\phi(0) = \phi_0, \quad (1b)$$

$$\partial_t \phi(0) = \phi'_0, \quad (1c)$$

where  $\phi: [0, T] \rightarrow X$ ,  $L$  is a closed, densely defined, self-adjoint, non-negative operator on  $X$  with domain  $\mathcal{D}(L)$ ,  $c$  is a positive constant, and  $f: \mathcal{D}(L) \times [0, T] \rightarrow X$  a smooth function. Such equations arise, for example, in acoustics, electromagnetics, quantum mechanics, and geophysical fluid dynamics, both in the “relativistic” ( $c = 1$ ) and “nonrelativistic” ( $c \gg 1$ ) regimes. The motivation for studying (1) as written is that it covers two well-studied special cases:

- (i)  $X = H^r(\mathbb{T}^d)$ ,  $L = \Delta$ , and  $f(\phi, t) = |\phi|^2 \phi$ , which corresponds to the standard semilinear Klein–Gordon equation.
- (ii)  $X = \mathbb{R}^{2d}$ ,  $L = 0$ , and  $f(\phi, t) = -2e^{-c^2 t J} \nabla V(e^{c^2 t J} \phi)$ , where  $J$  is the canonical symplectic matrix in  $2d$  dimensions and  $V$  is a smooth potential. Changing variables via  $q = e^{c^2 t J} \phi$ , we can write the system in the standard form

$$\dot{q} = p, \quad (2a)$$

$$\frac{1}{2c^2} \dot{p} = Jp - \nabla V(q). \quad (2b)$$

This system has been used as a finite-dimensional toy model for rotating fluid flow, where the limit  $c \rightarrow \infty$  corresponds to a rapidly rotating earth. Analytically, the non-relativistic limit regime is well-studied for these two examples. We refer the reader to [10, 15] for the case of the Klein–Gordon equation and to

[9, 12, 13] for the case of system (2). Numerically, equation (1) is extensively studied in the relativistic regime [11, 18]. However, due to the high oscillatory character of the solutions when  $c$  is large, most numerical methods suffer from severe time step restriction in the non-relativistic regime.

Several authors have considered the problem of finding “asymptotics-preserving numerical schemes”, i.e., schemes that perform uniformly in this singular limit. Some of these schemes [10] are based on a modulated Fourier expansion of the exact solution [8, 14] where the highly oscillatory problem in (2) is reduced to a non-oscillatory limit Schrödinger equation for which no time step restriction is needed. Other schemes are based on multiscale expansions of the exact solution [3, 5]. Chartier *et al.* [6] recently introduced a new method which employs an averaging transformation to soften the stiffness of the problem, hence allowing standard schemes to retain their order of convergence. Baumstark *et al.* [4] construct first and second order uniformly accurate integrators for the Klein–Gordon equation with *cubic* nonlinearity by integrating the trigonometric products arising from a suitable mild formulation explicitly.

In this paper, we develop a family of high-order asymptotics-preserving schemes for (1) that do not require pre-computations tied to the specific nonlinearity  $f$  and have no restrictions in time step size. The construction of the new schemes is explained in Section 4. We outline here their main ingredients, where the first two follow the prior work [4]:

- (i) Reformulate (1) as a coupled first order system using a linear transformation.
- (ii) Factor out the rapidly rotating phase to make it explicit.
- (iii) Iterate the resulting mild formulation up to the desired order for the coupled first order system in the new variables.
- (iv) Use the quadrature rule developed in [16] to handle the high oscillatory integral in the resulting mild formulation and complete the construction of the scheme.

The remainder of the paper is structured as follows. In Section 2, we state some properties of the operator  $L$  within the framework of its functional calculus. In Section 3, we introduce quadrature rules for the approximation of highly oscillatory Banach-space-valued functions in specific settings that will fit the construction of our schemes. Section 4 is devoted to the detailed construction of the schemes. Our main result on the order of convergence of the schemes is stated and proved in Section 5.

## 2. PRELIMINARIES

We define the operators  $B_c = c^{-1} \sqrt{L + c^2}$  and  $A_c = c^2 B_c - c^2$ . These operators are well-defined via the spectral theorem for densely defined normal operators (e.g. [17]). Indeed, for any densely defined normal operator  $P: \mathcal{D}(P) \subseteq X \rightarrow X$ , there exists a unique spectral measure  $E_P$  on the Borel  $\sigma$ -algebra  $\mathcal{B}(\mathbb{C})$  into the orthogonal projections on  $X$  such that

$$P = \int_{\mathbb{C}} \lambda dE_P(\lambda) = \int_{\sigma(P)} \lambda dE_P(\lambda). \quad (3)$$

This integral representation of  $P$  allows us to define the assignments  $P \mapsto f(P)$  for any  $E$ -a.e. finite measurable function  $f$  by the formula

$$f(P) = \int_{\sigma(P)} f(\lambda) dE_P(\lambda) \quad (4)$$

with domain

$$\mathcal{D}(f(P)) = \left\{ x \in X : \int_{\sigma(P)} |f(\lambda)|^2 d\langle E_P(\lambda)x, x \rangle < \infty \right\}. \quad (5)$$

*Definition 1.* Let  $A, B$  be two densely defined normal operators. If  $\mathcal{D}(AB) \subseteq \mathcal{D}(BA)$  and  $AB = BA$  on  $\mathcal{D}(AB)$ , we write  $AB \subseteq BA$  and say that “ $A$  commutes with  $B$ .”

We fix in what follows an operator  $J \in \mathcal{L}(X)$  such that

$$JL \subseteq LJ, \quad J^* = -J, \quad \text{and} \quad J^2 = -I. \quad (6)$$

We now collect important elementary properties of the operators  $J, A_c,$  and  $B_c$ .

**Lemma 2.** *The operators  $J, A_c,$  and  $B_c$  satisfy the following properties.*

- (i)  $\mathcal{D}(L) \subseteq \mathcal{D}(A_c) = \mathcal{D}(B_c) = \mathcal{D}(JA_c) = \mathcal{D}(JB_c)$ ,
- (ii)  $\|A_c u\|_{\mathcal{D}(L^j)} \leq \frac{1}{2} \|u\|_{\mathcal{D}(L^{j+1})}$  for any  $j \in \mathbb{N}$ ,
- (iii)  $J$  and  $e^{tJ}$  commute with  $f(L)$  for any measurable function  $f: \mathbb{R} \rightarrow \mathbb{R}$  and  $t \in \mathbb{R}$ ; in particular,  $J$  and  $e^{tJ}$  commute with  $A_c, B_c,$  and  $B_c^{-1}$ ,
- (iv)  $e^{tJA_c}$  commutes with  $J$  and  $f(L)$  for any measurable function  $f: \mathbb{R} \rightarrow \mathbb{R}$  and  $t \in \mathbb{R}$ ,
- (v)  $\|e^{tJA_c}\| \leq 1$ , and
- (vi)  $\|(e^{tJA_c} - I)u\| \leq \frac{1}{2} |t| \|u\|_{\mathcal{D}(L)}$ .

*Proof.* The inclusion in (i) follows directly from

$$\int_{\sigma(L)} |\lambda + c^2| d\langle E_L(\lambda)u, u \rangle \leq (c^2 + \frac{1}{2}) \|u\|^2 + \frac{1}{2} \int_{\sigma(L)} |\lambda|^2 d\langle E_L(\lambda)u, u \rangle; \quad (7)$$

the remaining identities are obvious. To prove (ii), we note that, for  $\lambda \geq 0$

$$c\sqrt{\lambda + c^2} - c^2 \leq \frac{\lambda}{2}, \quad (8)$$

and

$$\|A_c u\|_{\mathcal{D}(L^j)}^2 = \int_{\sigma(L)} \left| c\sqrt{\lambda + c^2} - c^2 \right|^2 (1 + |\lambda|^2)^j d\langle E_L(\lambda)u, u \rangle. \quad (9)$$

For (iii), we recall that  $J$  is bounded and commutes with  $L$ . Thus, by [17, Proposition 5.15],  $J E_L(K) = E_L(K) J$  for all  $K \in \mathcal{B}(\mathbb{C})$ . Consequently,  $e^{tJ} E_L(K) = E_L(K) e^{tJ}$  for all  $K \in \mathcal{B}(\mathbb{C})$  and  $t \in \mathbb{R}$ . Then the claim is a direct consequence of [17, Proposition 4.23]. For (iv), note that

$$e^{tJA_c} = \int_{\mathbb{C}^2} e^{ct\lambda(\sqrt{\mu^2 + c^2} - c)} dE_J(\lambda) dE_L(\mu) \quad (10)$$

where the integral is with respect to the product measure  $E_J \otimes E_L(K_1 \times K_2) = E_J(K_1) E_L(K_2)$  for all  $K_1, K_2 \in \mathcal{B}(\mathbb{C})$ . Hence,  $E_J$  and  $E_L$  commute with  $E_J \otimes E_L$  in the sense that  $E_J \otimes E_L(K_1 \times K_2) E(K_3) = E(K_3) E_J \otimes E_L(K_1 \times K_2)$  for all  $K_1, K_2, K_3 \in \mathcal{B}(\mathbb{C})$ . Then, Once again, the claim follows from [17, Proposition 4.23]. Estimate (v) is a direct consequence of the skew-symmetry of  $J$ . Finally, to prove

estimate (vi), let  $u \in \mathcal{D}(L)$ . Since the spectrum of  $J$  is purely imaginary and  $|e^{ix} - 1|^2 \leq x^2$  for  $x \in \mathbb{R}$ , we estimate

$$\begin{aligned} \|(e^{tJA_c} - I)u\|^2 &= \int_{\sigma(J) \times \sigma(L)} |e^{tc\lambda(\sqrt{\mu^2+c^2}-c)} - 1|^2 \langle dE_J(\lambda) dE_L(\mu)u, u \rangle \\ &\leq t^2 \int_{\sigma(J) \times \sigma(L)} |c\lambda(\sqrt{\mu^2+c^2}-c)|^2 \langle dE_J(\lambda) dE_L(\mu)u, u \rangle \\ &\leq t^2 \|JA_c u\|^2. \end{aligned} \quad (11)$$

The claim then follows by estimate (ii).  $\square$

*Remark 3.* Lemma 2(iii) and (iv) imply that if  $P$  and  $Q$  are two operators such that  $P$  is bounded and  $PQ \subseteq QP$ , then  $\mathcal{D}(PQ) = \mathcal{D}(Q)$  and  $P(\mathcal{D}(Q)) \subseteq \mathcal{D}(Q)$ . In other words, the domain of  $Q$  is invariant under any *bounded* operator commuting with  $Q$ . In this paper, the analysis of the numerical schemes assumes solutions of (1) in  $\mathcal{D}(L)$  which is, in view of this remark, invariant under any bounded operator commuting with  $L$ , in particular  $J$ ,  $e^{tJ}$ , and  $e^{tJA_c}$ .

### 3. QUADRATURE FOR BANACH-SPACE-VALUED FUNCTIONS

In this section, let  $(X, \|\cdot\|)$  be a complex Banach space and  $\Omega \subset \mathbb{C}$  be open. A function  $F: \Omega \rightarrow X$  is analytic if it is differentiable, i.e., provided for every  $z_0 \in \Omega$  there exists  $F'(z_0) \in X$  such that

$$F'(z_0) = \lim_{z \rightarrow z_0} \frac{F(z) - F(z_0)}{z - z_0}. \quad (12)$$

The following simple lemma shows that estimates on the quadrature error for differentiable complex-valued functions directly imply a corresponding estimate for  $X$ -valued functions.

**Lemma 4.** *Let  $I$  be an open interval on the real line and  $\mu$  a measure on  $I$ , possibly discrete. Suppose a quadrature rule with nodes  $x_k \in I$  and weights  $\omega_k$ ,  $k = 1, \dots, N$  satisfies the error estimate*

$$\left| \int_I f(x) d\mu(x) - \sum_{k=1}^N \omega_k f(x_k) \right| \leq C(N, I) \sup_{x \in I} |f^{(p)}(x)|, \quad (13)$$

for some  $p \in \mathbb{N}$  and every  $f \in C^p(I, \mathbb{C})$ . Then the quadrature rule satisfies the error estimate

$$\left\| \int_I F(x) d\mu(x) - \sum_{k=1}^N \omega_k F(x_k) \right\| \leq C(N, I) \sup_{x \in I} \|F^{(p)}(x)\|, \quad (14)$$

where the integral is understood in the Bochner-sense, for every  $F \in C^p(I, X)$ .

*Proof.* Fix  $\psi \in X^*$ . Let

$$e_N = \int_I F(x) d\mu(x) - \sum_{k=1}^N \omega_k F(x_k). \quad (15)$$

Due to the properties of the Bochner integral,

$$\psi(e_N) = \int_a^b \psi \circ F(x) d\mu - \sum_{k=1}^N \omega_k \psi \circ F(x_k), \quad (16)$$

so that, applying (13) to  $f = \psi \circ F$ , we obtain

$$\begin{aligned} |\psi(e_N)| &\leq C(N, I) \sup_{x \in I} \left| \frac{d^p}{dx^p} [\psi \circ F](x) \right| \\ &\leq C(N, I) \|\psi\|_* \sup_{x \in I} \|F^{(p)}(x)\|. \end{aligned} \quad (17)$$

By the Hahn–Banach theorem, we can choose  $\psi \in X^*$  with  $\|\psi\|_* \leq 1$  such that  $\psi(e_N) = \|e_N\|$ . This implies (14).  $\square$

With the help of this lemma, we lift three known estimates for the quadrature error of complex-valued functions to the Banach space setting. The first concerns the trapezoidal rule approximation for the integral of a 1-periodic function  $F$ , namely the uniformly weighted Riemann sum

$$T_N(F) = \frac{1}{N} \sum_{k=0}^{N-1} F\left(\frac{k}{N}\right). \quad (18)$$

For given  $a > 0$ , let

$$\Omega_a = \{z \in \mathbb{C}: -a < \operatorname{Im} z < a\}. \quad (19)$$

Then the following estimate, proved for  $X = \mathbb{C}$  in [19], holds true.

**Theorem 5.** *Let  $F$  be an  $X$ -valued function, 1-periodic on the real line, analytic with  $\|F(z)\| \leq A$  on the strip  $\Omega_a$  for some  $a > 0$ . Then for any  $N \in \mathbb{N}$ ,*

$$\left\| \int_0^1 F(x) dx - T_N(F) \right\| \leq \frac{2A}{e^{aN} - 1}. \quad (20)$$

*The constant 2 is as small as possible.*

The second concerns the Gauss formula for the integral of a function  $F$  defined on the interval  $[-1, 1]$ ,

$$G_N(F) = \sum_{k=1}^N \omega_k F(\xi_k), \quad (21)$$

where the  $\xi_k$  are the zeros of the Legendre polynomial  $p_N$  of degree  $N$  and the weights are given by

$$\omega_k = \frac{2}{(1 - \xi_k^2) [p'_N(\xi_k)]^2}. \quad (22)$$

For given  $b > a$  and  $\rho > \frac{1}{2}(b - a)$ , let  $E_\rho(a, b)$  denote the ellipse with foci  $a, b$  such that the lengths of its minor and major semiaxes sum up to  $\rho$ . Namely,

$$E_\rho(a, b) = \left\{ z \in \mathbb{C}: z = \frac{1}{2}(\rho e^{i\theta} + \frac{1}{4}(b - a)^2 \rho^{-1} e^{-i\theta}) + \frac{1}{2}(a + b), 0 \leq \theta < 2\pi \right\}, \quad (23)$$

and  $\Sigma_\rho(a, b)$  the open region in  $\mathbb{C}$  bounded by  $E_\rho(a, b)$ .

The formula (21) can easily be written out for functions defined on an arbitrary interval  $[a, b]$  using the affine change of variables

$$\ell: \Sigma_{\frac{2\rho}{b-a}}(-1, 1) \rightarrow \Sigma_\rho(a, b), \quad \ell(x) = \frac{b-a}{2}(x+1) + a. \quad (24)$$

**Theorem 6.** Fix  $M \in \mathbb{N}$ ,  $\varepsilon > 0$ , and  $\rho > \frac{1}{2}(b-a)$ . Let  $F: [a\varepsilon, b\varepsilon] \times \Sigma_\rho(a, b) \rightarrow X$  be such that  $x \mapsto F(\varepsilon x, z)$  is  $M$ -times differentiable for any  $z \in \Sigma_\rho(a, b)$  and  $\partial_x^k F(\varepsilon a, \cdot)$  is analytic on  $\Sigma_\rho(a, b)$  for  $k = 0, \dots, M-1$  with

$$\max_k \sup_{z \in \Sigma_\rho(a, b)} \|\partial_x^k F(\varepsilon a, z)\| \leq A_{\text{an}}, \quad (25a)$$

$$\sup_{\xi, x \in [a, b]} \|\partial_x^M F(\xi, x)\| \leq A_{\text{dif}}. \quad (25b)$$

We abbreviate  $f(x) = F(\varepsilon x, x)$ . Then, for any  $N \in \mathbb{N}$ ,

$$\left\| \int_a^b f(x) dx - G_N(f) \right\| \leq \frac{16 A_{\text{an}} e^{\varepsilon(b-a)} \rho^2}{(2\rho - b + a)} \left( \frac{b-a}{2\rho} \right)^{2N+1} + \frac{2 A_{\text{dif}} (b-a)^{M+1} \varepsilon^M}{M!}. \quad (26)$$

The formula  $G_N(f)$  is defined with nodes  $\eta_k = \ell(\xi_k)$ .

*Proof.* Writing the Taylor series with respect to the first variable of  $F$ , we find that for every  $x \in [a, b]$  there exists  $\xi = \xi(x) \in [a, b]$  such that

$$f(x) = \sum_{k=0}^{M-1} \frac{(x-a)^k \varepsilon^k}{k!} \partial_x^k F(\varepsilon a, x) + \frac{(x-a)^M \varepsilon^M}{M!} \partial_x^M F(\varepsilon \xi, x). \quad (27)$$

Thus, the following estimate, proved for  $X = \mathbb{C}$  in [7], holds true for the quadrature formula (21) applied on each  $f_k(z) = (z-a)^k \partial_x^k F(\varepsilon a, z)$ ,  $k = 0, \dots, M-1$ , which is analytic and bounded on  $\Sigma_\rho(a, b)$ :

$$\begin{aligned} & \left\| \int_a^b f_k(x) dx - G_N(f_k) \right\| \\ & \leq \frac{16 \rho^2}{(2\rho - b + a)} \left( \frac{b-a}{2\rho} \right)^{2N+1} \sup_{z \in \Sigma_\rho(a, b)} \|(z-a)^k \partial_x^k F(\varepsilon a, z)\|. \end{aligned} \quad (28)$$

This yields the first term on the right of (26). The Lagrange remainder in (27) is estimated independently for the continuum integral over the interval  $[a, b]$  and for the discrete integral  $G_N$ , in both cases yielding a contribution to the second term on the right of (26).  $\square$

The third concerns the Gauss formula for the discrete sum  $\sum_{j=0}^{m-1} F(x_j)$  on equidistant nodes

$$x_j = -1 + \frac{2j}{m-1}, \quad 0 \leq j \leq m-1 \quad (29)$$

with

$$\frac{2}{m} \sum_{j=0}^{m-1} F(x_j) \approx S_M(F) \equiv \sum_{k=1}^M \omega_{k,m} F(\xi_{k,m}), \quad (30)$$

where the quadrature nodes  $\xi_{k,m}$  are the zeros of the so-called Gram polynomial  $p_M$  of degree  $M$ . Such polynomials are defined by their orthonormality with respect to the discrete equidistant sum, namely

$$\sum_{j=0}^{m-1} p_{n,m}(x_j) p_{k,m}(x_j) = \delta_{nk}. \quad (31)$$

The weights  $\omega_{k,m}$  are given by

$$\omega_{k,m} = \frac{a_{M,m}}{a_{M-1,m}} \frac{2}{m p'_{M,m}(\xi_{k,m}) p_{M-1,m}(\xi_{k,m})}, \quad (32)$$

where  $a_{M,m}$  denotes the leading coefficient of  $p_{M,m}$ . For a detailed derivation and discussion, see [1, 2, 16].

**Theorem 7.** Fix  $M \in \mathbb{N}$  and let  $F: [a, b] \rightarrow X$  be a  $2M$ -times differentiable function with  $\|F^{(2M)}\| \leq A$  on  $[a, b]$ . Then

$$\left\| \frac{b-a}{m-1} \sum_{j=0}^{m-1} F(y_j) - \frac{m(b-a)}{2(m-1)} S_M(F) \right\| \leq \frac{16 A (b-a)^{2M+1} M!^4}{(2M+1)(2M)!^3}. \quad (33)$$

The formula  $S_M(f)$  is defined with nodes  $\eta_{k,m} = \ell(\xi_{k,m})$  and the equidistant summation points are given by  $y_j = \ell(x_j)$ .

*Proof.* Assume first that  $a = -1$  and  $b = 1$ . The general case then follows via the affine change of variable  $\ell(x) = \frac{b-a}{2}(x+1) + a$ . Assume further that  $X = \mathbb{R}$ . The general case where  $X$  is a complex Banach space follows by applying Lemma 4.

Thus, let  $H$  be the unique polynomial of degree  $2M-1$  satisfying the Hermite interpolation problem

$$F(\xi_{k,m}) = H(\xi_{k,m}), \quad F'(\xi_{k,m}) = H'(\xi_{k,m}), \quad k = 1, \dots, M. \quad (34)$$

By Rolle's theorem, for any  $x \in [-1, 1]$  there exists  $\xi(x) \in [-1, 1]$  such that

$$F(x) - H(x) = \frac{F^{(2M)}(\xi)}{(2M)!} q_{M,m}^2(x), \quad (35)$$

where  $q_{M,m}$  is the polynomial

$$q_{M,m}(x) = \prod_{k=1}^M (x - x_{k,m}). \quad (36)$$

Since (30) is exact for all polynomials of degree less than  $2M-1$ ,

$$\frac{2}{m} \sum_{j=0}^{m-1} F(x_j) = S_M(H) = S_M(F). \quad (37)$$

Thus, using (35), we estimate

$$\begin{aligned} \left\| \frac{2}{m-1} \sum_{j=0}^{m-1} F(x_j) - \frac{m}{m-1} S_M(F) \right\| &= \frac{2}{m} \sum_{j=0}^{m-1} \|F(x_j) - H(x_j)\| \\ &\leq \frac{2A}{(m-1)(2M)!} \sum_{j=0}^{m-1} q_{M,m}^2(x_j). \end{aligned} \quad (38)$$

Note that  $\deg(q_{M,m}) = M$  and  $q_{M,m}$  has the same zeros as the Gram polynomial  $p_{M,m}$ . Hence,

$$p_{M,m} = a_{M,m} q_{M,m}, \quad (39)$$

where the constant  $a_{M,m}$  is given by [16]

$$a_{M,m} = \sqrt{\frac{(2m+1)(m-M-1)! (2M)! (m-1)^M}{(m+M)! 2^M M!^2}}. \quad (40)$$

Since  $p_{M,m}$  is normalized,

$$\sum_{j=0}^{m-1} q_{M,m}^2(x_j) = \frac{1}{a_{M,m}^2} = \frac{(m+M)!}{(2M+1)(m-M-1)!(m-1)^{2M}} \frac{2^M M!^4}{(2M)!^2}. \quad (41)$$

For  $m > M \geq 1$ , we have

$$\begin{aligned} \frac{(m+M)!}{(m-M-1)!(m-1)^{2M+1}} &= \frac{(m+M)(m+M-1)\cdots(m-M)}{(m-1)(m-1)\cdots(m-1)} \\ &\leq \left(1 + \frac{1}{M}\right)^{2M+1} \\ &\leq 8, \end{aligned}$$

which completes the proof.  $\square$

In the next section, we will need to approximate a double integral of a function of two variables where one of the integrals is continuous, the other discrete. The corresponding estimate, combining (33) and (26) is as follows.

**Lemma 8.** *Fix  $0 < \tau < 1$ ,  $m \in \mathbb{N}$ , and set  $T = \tau/m$ . Suppose that the function  $G: [0, \tau] \times [0, 1] \rightarrow X$  satisfies the following.*

(i) *For any  $x \in [0, 1]$ ,  $s \mapsto G(s, x)$  is  $2M$ -times differentiable with*

$$\sup_{[0, \tau] \times [0, 1]} \|\partial_s^{2M} G\| \leq A. \quad (42)$$

(ii)  *$G(s, x) = F(s, Tx, x)$  for some function  $F: [0, \tau] \times [0, T] \times \Sigma_{1/(2\gamma)}(0, 1) \rightarrow X$  with  $\gamma \in (0, 1)$  such that for every  $s \in [0, \tau]$ ,  $F(s, \cdot, \cdot)$  satisfies the assumptions of Theorem 6 replacing  $M$  there by  $2M$  here, and with bounds that are uniform with respect to  $s \in [0, \tau]$ .*

Then there exists a constant  $C > 0$  depending on  $G$  such that for any  $N$ ,

$$\left\| T \sum_{j=0}^{m-1} \int_0^1 G(jT, x) dx - \frac{\tau}{4} \sum_{i=1}^M \sum_{k=1}^N \omega_{i,m} \omega_k G(\eta_{i,m}, \eta_k) \right\| \leq C \tau (\gamma^{2N} + \tau^{2M}). \quad (43)$$

*Proof.* Using Theorem 7, we find that

$$T \sum_{j=0}^{m-1} G(jT, x) = \frac{\tau}{2} \sum_{i=1}^M \omega_{i,m} G(\eta_{i,m}, x) + R(x), \quad (44)$$

where, in view of the first assumption on  $G$  and by using estimate (33), there exists a constant  $C$  depending on  $\sup_{[0, \tau] \times [0, 1]} \|\partial_s^{2M} G\|$  such that

$$\|R(x)\| \leq C \tau^{2M+1}. \quad (45)$$

For each node  $\eta_{i,m}$ , the function  $F(\eta_{i,m}, \cdot, \cdot)$  satisfies the assumption of Theorem 6 on  $[0, T] \times \Sigma_{1/(2\gamma)}(0, 1)$ . Thus, taking the integral over  $[0, 1]$  of (44) and using estimate (26), we obtain (43).  $\square$

## 4. UNIFORMLY ACCURATE SCHEMES

Following [4], we introduce “twisted variables” in which the linear operator in the equation is uniform as  $c \rightarrow \infty$ . In a first change of variables, we set

$$U = \phi - c^{-2} B_c^{-1} J \dot{\phi}, \quad (46a)$$

$$V = \phi + c^{-2} B_c^{-1} J \dot{\phi}. \quad (46b)$$

In terms of the variables  $U$  and  $V$ , equation (1) reads

$$J \dot{U} = -c^2 B_c U + B_c^{-1} f\left(\frac{1}{2}(U + V), t\right), \quad (47a)$$

$$J \dot{V} = c^2 B_c V - B_c^{-1} f\left(\frac{1}{2}(U + V), t\right). \quad (47b)$$

As a second change of variables, we define

$$u = e^{-c^2 t J} U, \quad v = e^{c^2 t J} V. \quad (48)$$

In terms of  $u$  and  $v$ , system (47) takes the form

$$\dot{u} = J A_c u - J B_c^{-1} e^{-c^2 t J} f\left(\frac{1}{2}(e^{c^2 t J} u + e^{-c^2 t J} v), t\right), \quad (49a)$$

$$\dot{v} = -J A_c v + J B_c^{-1} e^{c^2 t J} f\left(\frac{1}{2}(e^{c^2 t J} u + e^{-c^2 t J} v), t\right). \quad (49b)$$

We can write this system more compactly in terms of the vector-valued functions  $W = (U, V)^\top$  and  $w = (u, v)^\top$ . Letting  $A_c$  and  $B_c$  act diagonally on  $\mathcal{D}(A_c) \times \mathcal{D}(A_c)$  and defining

$$\mathcal{J} = \begin{pmatrix} J & 0 \\ 0 & -J \end{pmatrix}, \quad (50a)$$

$$\mathcal{F}(W, t) = (-J, J)^\top f\left(\frac{1}{2}(U + V), t\right), \quad (50b)$$

we can write

$$\dot{w} = \mathcal{J} A_c w + B_c^{-1} e^{-c^2 t \mathcal{J}} \mathcal{F}(e^{c^2 t \mathcal{J}} w). \quad (51)$$

Applying the Duhamel formula, we obtain

$$\begin{aligned} w(t_n + \tau) &= e^{\tau \mathcal{J} A_c} w(t_n) \\ &\quad - B_c^{-1} \int_0^\tau e^{(\tau-s) \mathcal{J} A_c} e^{-c^2(t_n+s) \mathcal{J}} \mathcal{F}(e^{c^2(t_n+s) \mathcal{J}} w(t_n + s), s) ds. \end{aligned} \quad (52)$$

The two following assumptions on the nonlinearity  $f$  and on the solution of (52) are required for the rigorous construction of the numerical scheme.

*Assumption 9.* For given  $M \in \mathbb{N}$ , we assume that  $f$  satisfies the following

- (i)  $t \mapsto f(u, t)$  is analytic for fixed  $u \in \mathcal{D}(L^M)$  and  $x \mapsto f(e^{2\pi x J} u, t)$  has an analytic extension to  $\Sigma_{\frac{1}{2\gamma}}(0, 1)$  for some  $\gamma \in (0, 1)$  for fixed  $t \in [0, T]$ .
- (ii) For fixed  $t \in [0, T]$ ,  $f(\cdot, t): \mathcal{D}(L^M) \rightarrow X$  is  $M$ -times Gâteaux differentiable such that  $D^k f(u, t) \in \mathcal{L}(\mathcal{D}(L^{M-\alpha_k}), \mathcal{D}(L^{M-|\alpha_k|}))$  for every  $k = 1, \dots, M$ ,  $u \in \mathcal{D}(L^M)$ , and multi-index  $\alpha_k = (j_1, \dots, j_k)$  for which each component is larger than 1 and  $|\alpha_k| \leq M$ .

Here,  $\mathcal{D}(L^{M-\alpha_k})$  refers to the direct product  $\mathcal{D}(L^{M-j_1}) \times \dots \times \mathcal{D}(L^{M-j_k})$ .

*Remark 10.* The nonlinearity of the semilinear Klein–Gordon equation introduced in Section 1 satisfies Assumption 9.

*Remark 11.* To see how condition (ii) in Assumption 9 arises, consider the following example, which is a simplified version of the estimates which arise in the analysis of the numerical scheme below. Take  $g(x) = e^{xJA_c} f(h(x))$ ,  $h(x) = e^{xJA_c} u$ ,  $u \in \mathcal{D}(L^M)$  and  $M = 2$ . Since

$$\begin{aligned} e^{-xJA_c} g''(x) &= -A_c^2 f(h(x)) + 2JA_c Df(h(x)) h'(x) \\ &\quad + D^2 f(h(x)) [h'(x), h'(x)] + Df(h(x)) h''(x), \end{aligned} \quad (53)$$

$\|g''\|_X$  is uniformly bounded in  $c$  provided

$$Df(u) \in \mathcal{L}(\mathcal{D}(L^{M-1}), D(L^{M-1})) = \mathcal{L}(\mathcal{D}(L)), \quad (54a)$$

$$Df(u) \in \mathcal{L}(\mathcal{D}(L^{M-2}), D(L^{M-2})) = \mathcal{L}(X), \quad (54b)$$

$$D^2 f(u) \in \mathcal{L}(\mathcal{D}(L^{M-(1,1)}), \mathcal{D}(L^{M-2})) = \mathcal{L}(\mathcal{D}(L) \times \mathcal{D}(L), X). \quad (54c)$$

This suffices to then satisfy condition (ii) of Lemma 2 for  $G = g$ .

*Assumption 12.* For given  $M$ , there exists  $\mathcal{T} > 0$  and  $K > 0$  independent of  $c$  such that

$$\sup_{0 \leq t \leq \mathcal{T}} \|w(t)\|_{\mathcal{D}(L^M)} \leq K. \quad (55)$$

To guarantee uniform convergence with respect to  $c$ , we make the following important observation which effectively asserts that the time derivative  $\dot{w}$  is bounded uniformly in  $c$ .

**Lemma 13.** *The solution  $w$  of (52) satisfies*

$$\|w(t_n + s) - w(t_n)\| \leq \frac{s}{2} \|w(t_n)\|_{\mathcal{D}(L)} + s \sup_{\xi \in [0, s]} \|\mathcal{F}(e^{c^2(t_n + \xi)\mathcal{J}} w(t_n + \xi))\|. \quad (56)$$

*Proof.* The proof is a direct application of estimate (vi) in Lemma 2 and the fact that  $\|B_c^{-1}\| \leq 1$ .  $\square$

Since we can adapt the time step  $\tau$  of what is to emerge as the numerical scheme, it is most convenient to select  $\tau$  as an integer multiple of the fast period  $T = 2\pi/c^2$  so that  $\tau = mT$  for some  $m \in \mathbb{N}$ . As  $e^{s\mathcal{J}} = \cos(s)I + \sin(s)\mathcal{J}$ ,

$$e^{jc^2T\mathcal{J}} = e^{2\pi j\mathcal{J}} = I \quad (57)$$

whenever  $j$  is integer.

In a first step, we define a sequence of “pre-schemes”  $\Phi_l: X \times \mathbb{R} \rightarrow X$  which provide consistent approximations to the right hand side of the Duhamel formula (52) to order  $\tau^{l+1}$ , namely

$$\Phi_1(w, z) = e^{z\mathcal{J}A_c} w - B_c^{-1} \int_0^z e^{-c^2 s\mathcal{J}} \mathcal{F}(e^{c^2 s\mathcal{J}} w, s) ds, \quad (58a)$$

$$\Phi_{l+1}(w, z) = e^{z\mathcal{J}A_c} w - B_c^{-1} \int_0^z e^{(z-s)\mathcal{J}A_c} e^{-c^2 s\mathcal{J}} \mathcal{F}(e^{c^2 s\mathcal{J}} \Phi_l(w, s), s) ds. \quad (58b)$$

The pre-schemes approximate the true solution in the following sense.

**Lemma 14.** *Let  $f$  be Lipschitz on  $X$ , let  $w$  be a solution for (52) satisfying Assumption 12 for  $M = 1$ , and fix  $l \in \mathbb{N}^*$ . Then there exist constants  $C_l$  independent of  $c$  such that all  $s \geq 0$ ,*

$$\|w(t_n + s) - \Phi_l(w(t_n), s)\| \leq C_l s^{l+1}. \quad (59)$$

*Proof.* We set  $R_l(w(t_n), s) = w(t_n + s) - \Phi_l(w(t_n), s)$  and proceed by induction. When  $l = 1$ ,

$$\begin{aligned} R_1(w(t_n), s) &= B_c^{-1} \int_0^s e^{-c^2 \sigma \mathcal{J}} \mathcal{F}(e^{c^2 \sigma \mathcal{J}} w(t_n), \sigma) d\sigma \\ &\quad - B_c^{-1} \int_0^s e^{(s-\sigma)\mathcal{J}A_c} e^{-c^2 \sigma \mathcal{J}} \mathcal{F}(e^{c^2 \sigma \mathcal{J}} w(t_n + \sigma), \sigma) d\sigma. \end{aligned} \quad (60)$$

The estimate on  $R_1$  follows by using Lemma 13 to freeze  $w(t_n + \sigma)$  and Lemma 2(iv) to remove the operator  $e^{(s-\sigma)\mathcal{J}A_c}$  in the second integral in (60). For  $l \geq 1$ ,

$$\begin{aligned} R_{l+1}(w(t_n), s) &= B_c^{-1} \int_0^s e^{(s-\sigma)\mathcal{J}A_c} e^{-c^2 \sigma \mathcal{J}} \mathcal{F}(e^{c^2 \sigma \mathcal{J}} \Phi_l(w(t_n), \sigma), \sigma) d\sigma \\ &\quad - B_c^{-1} \int_0^s e^{(s-\sigma)\mathcal{J}A_c} e^{-c^2 \sigma \mathcal{J}} \mathcal{F}(e^{c^2 \sigma \mathcal{J}} w(t_n + \sigma), \sigma) d\sigma. \end{aligned} \quad (61)$$

By Lemma 2 and the fact that  $f$  is Lipschitz on  $X$ , there exists a constant  $C$  independent of  $c$  such that

$$\|R_{l+1}(w(t_n), s)\| \leq C s \sup_{\sigma \leq s} \|R_l(w(t_n), \sigma)\|. \quad (62)$$

This completes the proof.  $\square$

While the operator  $A_c$  and the associated semi-group  $e^{t\mathcal{J}A_c}$  are uniformly well-behaved as  $c \rightarrow \infty$ , the integrals in (58) still contain highly oscillatory terms with a *single* fast frequency. For the latter, effective numerical quadrature is possible [16]. Following the strategy developed there, we split  $z/T \equiv m_z + \theta_z$  into its integer part  $m_z = \lfloor z/T \rfloor$  and fractional part  $\theta_z = z/T - m_z$ . Then the integral in (58a) can be written

$$\begin{aligned} &B_c^{-1} \int_0^z e^{-c^2 s \mathcal{J}} \mathcal{F}(e^{c^2 s \mathcal{J}} w, s) ds \\ &= B_c^{-1} \sum_{j=0}^{m_z-1} \int_{jT}^{(j+1)T} e^{-c^2 s \mathcal{J}} \mathcal{F}(e^{c^2 s \mathcal{J}} w, s) ds \\ &\quad + B_c^{-1} \int_{m_z T}^z e^{-c^2 s \mathcal{J}} \mathcal{F}(e^{c^2 s \mathcal{J}} w, s) ds \\ &= T \sum_{j=0}^{m_z-1} \int_0^1 G_0(jT, \sigma) d\sigma + T \int_0^{\theta_z} G_0(m_z T, \sigma) d\sigma \end{aligned} \quad (63)$$

with

$$G_0(\rho, \sigma) = B_c^{-1} e^{-2\pi\sigma\mathcal{J}} \mathcal{F}(e^{2\pi\sigma\mathcal{J}} w, \rho + \sigma T) \quad (64)$$

and where, in the second equality of (63), we have used (57). Analogously, the integral in (58b) can be written

$$\begin{aligned} &B_c^{-1} \int_0^z e^{-s\mathcal{J}A_c} e^{-c^2 s \mathcal{J}} \mathcal{F}(e^{c^2 s \mathcal{J}} \Phi_l(w, s), s) ds \\ &= T \sum_{j=0}^{m_z-1} \int_0^1 G[\Phi_l](jT, \sigma) d\sigma + T \int_0^{\theta_z} G[\Phi_l](m_z T, \sigma) d\sigma, \end{aligned} \quad (65)$$

where, for  $\Upsilon: X \times \mathbb{R} \rightarrow X$ ,

$$G[\Upsilon](\rho, \sigma) = B_c^{-1} e^{-(\rho+\sigma T)\mathcal{J}A_c} e^{-2\pi\sigma\mathcal{J}} \mathcal{F}(e^{2\pi\sigma\mathcal{J}} \Upsilon(w, \rho + \sigma T), \rho + \sigma T). \quad (66)$$

Altogether, (58) then takes the form

$$\Phi_1(w, z) = e^{z\mathcal{J}A_c} w - T \sum_{j=0}^{m_z-1} \int_0^1 G_0(jT, \sigma) d\sigma - T \int_0^{\theta_z} G_0(m_z T, \sigma) d\sigma, \quad (67a)$$

$$\Phi_{l+1}(w, z) = e^{z\mathcal{J}A_c} \left( w - T \sum_{j=0}^{m_z-1} \int_0^1 G[\Phi_l](jT, \sigma) d\sigma - T \int_0^{\theta_z} G[\Phi_l](m_z T, \sigma) d\sigma \right). \quad (67b)$$

We now use the approximation described in Lemma 8 to define a sequence of numerical schemes,

$$\begin{aligned} \Psi_1(w, z) &= e^{z\mathcal{J}A_c} w - \frac{m_z T}{4} \sum_{i=1}^M \sum_{k=0}^N \omega_{i, m_z} \omega_k G_0(\eta_{i, m_z}, \eta_k) \\ &\quad - \frac{\theta_z T}{2} \sum_{k=0}^N \omega_k G_0(m_z T, \theta_z \eta_k), \end{aligned} \quad (68a)$$

$$\begin{aligned} \Psi_{l+1}(w, z) &= e^{z\mathcal{J}A_c} \left( w - \frac{m_z T}{4} \sum_{i=1}^M \sum_{k=0}^N \omega_{i, m_z} \omega_k G[\Psi_l](\eta_{i, m_z}, \eta_k) \right. \\ &\quad \left. - \frac{\theta_z T}{2} \sum_{k=0}^N \omega_k G[\Psi_l](m_z T, \theta_z \eta_k) \right). \end{aligned} \quad (68b)$$

*Remark 15.* For a scheme of global order  $l$ , we use  $\Psi_l$  with  $z = \tau$  as the time stepper. At the top level, the second sum in (68a) or (68b) does not contribute. However, when  $l \geq 2$ , the inner evaluations of  $\Psi_{l-1}, \Psi_{l-2}, \dots$  will generally be evaluated at points  $z$  that are not integer multiples of  $T$ , so that their  $z$ -arguments have to be re-split into the respective integer ( $m_z$ ) and fractional ( $\theta_z$ ) multiples of  $T$ . Thus, in general, the second sum on the right of (68) is required for consistency.

*Remark 16.* Note that in the case where  $\mathcal{F}$  is constant with respect to the second variable, the function  $G_0 = G_0(x)$  is one-variable periodic function. Thus, the approximation from Theorem 5 can also be used to define a first order scheme ( $l = 1$ ) with accuracy that is exponential in the number of nodes. More specifically, for  $\tau = mT$ ,

$$\begin{aligned} \Phi_1(w, \tau) &= e^{\tau\mathcal{J}A_c} w - \tau \int_0^1 G_0(x) dx \\ &= e^{\tau\mathcal{J}A_c} w - \frac{\tau}{N} \sum_{k=0}^{N-1} G_0\left(\frac{k}{N}\right) + \mathcal{O}(\tau e^{-dN}) \end{aligned} \quad (69)$$

for some  $d > 0$ .

**Lemma 17.** *Let  $l, M \in \mathbb{N}^*$  and  $w \in \mathcal{D}(L^{2M})$ . Fix  $0 < \gamma < 1$  and assume that  $f$  satisfies Assumption 9 with  $M$  there replaced by  $2M$  here. Then there exists a constant  $C_l$  depending only on  $\gamma$  and  $\|w\|_{\mathcal{D}(L^{2M})}$  such that for all  $N \in \mathbb{N}^*$  and  $z < 1$ ,*

$$\|\Psi_l(w, z) - \Phi_l(w, z)\| \leq C_l z (z^{2M} + \gamma^{2N}). \quad (70)$$

*Proof.* Since  $w \in \mathcal{D}(L^{2M})$  and  $f$  satisfies Assumption 9 the two functions the functions  $G_0$  and  $G[\Phi_l]$  satisfy the conditions of Lemma 8 on  $[0, z] \times [0, 1] \times \Sigma_{1/(2\gamma)}(0, 1)$ .

We set  $S_l(w, z) = \Psi_l(w, z) - \Phi_l(w, z)$  and proceed by induction. For  $l = 1$ , we can directly use Lemma 8 for the difference of first terms and Theorem 6 for the difference of second terms, (70) holds true as stated. For  $l > 1$ , we have

$$\begin{aligned}
& e^{-z\mathcal{J}A_c} S_{l+1}(w, z) \\
&= -\frac{m_z T}{4} \sum_{i=1}^M \sum_{k=0}^N \omega_{i, m_z} \omega_k G[\Psi_l](\eta_{i, m_z}, \eta_k) - \frac{\theta_z T}{2} \sum_{k=0}^N \omega_k G[\Psi_l](m_z T, \theta_z \eta_k) \\
&+ \frac{m_z T}{4} \sum_{i=1}^M \sum_{k=0}^N \omega_{i, m_z} \omega_k G[\Phi_l](\eta_{i, m_z}, \eta_k) + \frac{\theta_z T}{2} \sum_{k=0}^N \omega_k G[\Phi_l](m_z T, \theta_z \eta_k) \\
&- \frac{m_z T}{4} \sum_{i=1}^M \sum_{k=0}^N \omega_{i, m_z} \omega_k G[\Phi_l](\eta_{i, m_z}, \eta_k) - \frac{\theta_z T}{2} \sum_{k=0}^N \omega_k G[\Phi_l](m_z T, \theta_z \eta_k) \\
&+ T \sum_{j=0}^{m_z-1} \int_0^1 G[\Phi_l](jT, \sigma) d\sigma + T \int_0^{\theta_z} G[\Phi_l](m_z T, \sigma) d\sigma \tag{71}
\end{aligned}$$

We write  $S_{l+1}^{(1)}(w, z)$  and  $S_{l+1}^{(2)}(w, z)$  to denote the first two and the last two lines on the right of (71), respectively. As  $\mathcal{F}$  is Lipschitz on  $X$ , there exist  $K_1, K_2$  independent of  $c$  such that

$$\begin{aligned}
\|S_{l+1}^{(1)}(w, z)\| &\leq K_1 z \sup_{i,k} \|S_l(w, \eta_{i, m_z} + \eta_k T)\| + K_2 z \sup_k \|S_l(w, m_z T + \theta_z \eta_k T)\| \\
&\leq (K_1 + K_2) C_l z^2 (z^{2M} + \gamma^{2N}). \tag{72}
\end{aligned}$$

On the other hand, using Lemma 8, there exists a constant  $D_2 = D_2(\|w\|_{\mathcal{D}(L^{2M})}, \gamma)$ , such that

$$\|S_{l+1}^{(2)}(w, z)\| \leq D_2 z (z^{2M} + \gamma^{2N}). \tag{73}$$

Thus, combining (72) and (73), we conclude that there exists  $C_{l+1}$  depending on  $\|w\|$  and  $\gamma$  such that

$$\|S_{l+1}(w, z)\| \leq C_{l+1} z (z^{2M} + \gamma^{2N}), \tag{74}$$

which concludes the proof.  $\square$

As stated before, we select the time step  $\tau$  to be an integer multiple of the fast period  $T$  so that  $\tau = mT$  for some  $m \in \mathbb{N}$ . As a numerical approximation to the exact solution  $w$  at time  $t_{n+1}$ , we take the scheme

$$w_{n+1} = \Psi_l(w_n, \tau), \tag{75a}$$

$$w_0 = \begin{pmatrix} \phi_0 \\ \phi_0 \end{pmatrix} - c^{-2} \mathcal{J} B_c^{-1} \begin{pmatrix} \phi_0' \\ \phi_0' \end{pmatrix}. \tag{75b}$$

## 5. CONVERGENCE ANALYSIS

The scheme (75) satisfies the following global estimate.

**Theorem 18.** *Fix  $l \in \mathbb{N}^*$ ,  $\gamma < 1$ ,  $c_0 > 0$ , and let  $M = \lfloor \frac{l+1}{2} \rfloor$ . Assume further that there exists  $\mathcal{K} > 0$  such that for every  $c \geq c_0$ ,*

$$\|\phi_0\|_{\mathcal{D}(L^{2M})} + c^{-2} \|B_c^{-1} \phi_0'\|_{\mathcal{D}(L^{2M})} \leq \mathcal{K}. \tag{76}$$

Then there exist  $\mathcal{T} > 0$  and  $C = C(\gamma, \mathcal{K}, \mathcal{T}, c_0)$  such that for all  $c \geq c_0$ ,  $\tau \in \frac{\pi}{c^2}\mathbb{N}$ ,  $t_n \leq \mathcal{T}$ , and  $N \in \mathbb{N}^*$ ,

$$\|\phi_n - \phi(t_n)\| \leq C(\tau^l + \gamma^{2N}) \quad (77)$$

where  $\phi$  solves (1) and

$$\phi_n = \frac{(w_n)_1 + (w_n)_2}{2} \quad (78)$$

with  $w_n$  given by (75).

*Proof.* Note first that for every  $c \geq c_0$ ,

$$\|w_0\|_{\mathcal{D}(L^{2M})} \leq \|\phi_0\|_{\mathcal{D}(L^{2M})} + \|c^{-2} JB_c^{-1} \phi'_0\|_{\mathcal{D}(L^{2M})} \leq \mathcal{K}. \quad (79)$$

Thus, there exist two constants  $\mathcal{T}, K > 0$  depending on  $c_0$  and  $w_0$  for which Assumption 12 is satisfied. Lemma 17 and 14 allow us to write

$$\begin{aligned} w(t_n + \tau) &= \Phi_l(w(t_n), \tau) + R_l(w(t_n), \tau) \\ &= \Psi_l(w(t_n), \tau) + R_l(w(t_n), \tau) - S_l(w(t_n), \tau). \end{aligned} \quad (80)$$

Setting  $e_n = \|w(t_n) - w_n\|$ , we now split the error as follows:

$$\begin{aligned} e_{n+1} &\leq \|R_l(w(t_n), \tau)\| + \|S_l(w(t_n), \tau)\| \\ &\quad + \|\Psi_l(w(t_n), \tau) - \Psi_l(w_n, \tau)\|. \end{aligned} \quad (81)$$

Recalling that  $\mathcal{F}$  is Lipschitz on  $X$  and arguing by induction on  $l$ , we find that there exists a constant  $C_1 > 0$  independent of  $c$  such that

$$\|\Psi_l(w(t_n), \tau) - \Psi_l(w_n, \tau)\| \leq (1 + C_1 \tau)^l e_n. \quad (82)$$

By Lemma 14 and 17, there exists a constant  $C_2 > 0$  depending on  $\gamma$  and  $K$  such that

$$\|R_l(w(t_n), \tau)\| + \|S_l(w(t_n), \tau)\| \leq C_2 \tau (\tau^l + \gamma^{2N}). \quad (83)$$

Then, (81) reads

$$e_{n+1} \leq (1 + C_1 \tau)^l e_n + C_2 \tau (\tau^l + \gamma^{2N}). \quad (84)$$

Thus, we find by induction that

$$e_n \leq (1 + C_1 \tau)^{n+l} e_0 + C_2 \frac{(1 + C_1 \tau)^{n+l} - 1}{C_1} (\tau^l + \gamma^{2N}). \quad (85)$$

Since  $e_0 = 0$  and  $1 + x \leq e^x$ , we obtain

$$e_n \leq C_2 \frac{e^{C_1 l \mathcal{T}} - 1}{C_1} (\tau^l + \gamma^{2N}) \equiv C (\tau^l + \gamma^{2N}). \quad (86)$$

To obtain the final estimate, we undo the variable twist, noting that

$$\phi(t_n) = \frac{(e^{c^2 t_n J}(w(t_n)))_1 + (e^{-c^2 t_n J}(w(t_n)))_2}{2} = \frac{w(t_n)_1 + w(t_n)_2}{2}. \quad (87)$$

Then (86) directly implies estimate (77).  $\square$

#### ACKNOWLEDGMENTS

The work was supported by German Research Foundation grant OL-155/6-2. MO further acknowledges support through German Research Foundation Collaborative Research Center TRR 181 under project number 274762653.

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