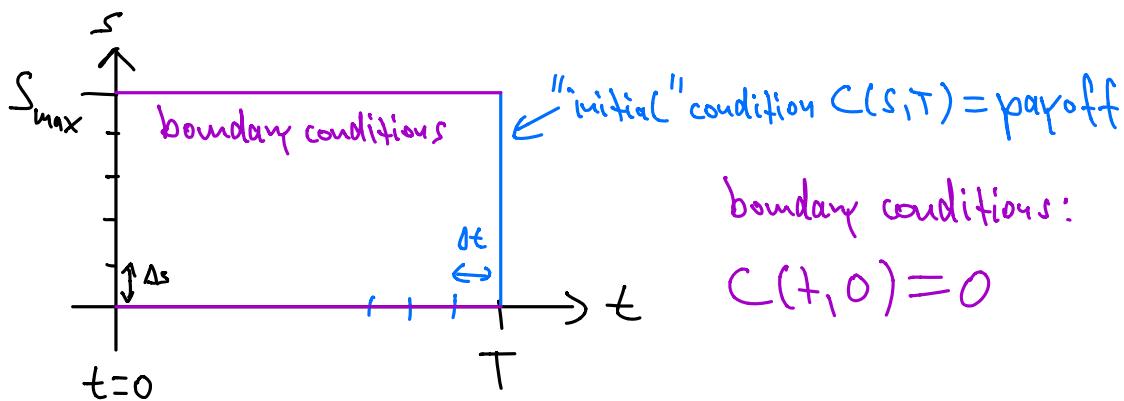


4.3 Discrete Finite Differences



boundary condition at S_{\max} is only due to discretization; good choices are:

- $C(t, S_{\max}) = S_{\max} - Ke^{-r(T-t)}$, this is a sol. to BS eq.
- $C(t, S_{\max}) = S_{\max}$, simpler; justified since S_{\max} should be chosen $\gg K$

Discretization:

- partition $[0, T]$ into M steps of size $\Delta t = \frac{T}{M}$, $t_j = j\Delta t$
- partition $[0, S_{\max}]$ into N steps of size $\Delta S = \frac{S_{\max}}{N}$, $s_i = i\Delta S$
- we abbreviate $C(t_j, s_i) = C_i^j$

Then:

$$\frac{\partial C_i^j}{\partial t} = \frac{C_i^{j+1} - C_i^j}{\Delta t} + O(\Delta t) \quad \text{for fixed } i$$

$$\frac{\partial C_i^j}{\partial S} = \frac{C_{i+1}^j - C_i^j}{\Delta S} + O(\Delta S) \quad \text{for fixed } j$$

the spatial derivative can be improved (fix j here):

Taylor

$$(1) \quad C(s_i + \Delta s) = C(s_i) + \frac{\partial C}{\partial s}(s_i) \Delta s + \frac{1}{2} \frac{\partial^2 C}{\partial s^2}(s_i) \Delta s^2 + \frac{1}{3!} \frac{\partial^3 C}{\partial s^3} \Delta s^3 + O(\Delta s^4)$$

$$(2) \quad C(s_i - \Delta s) = C(s_i) - \frac{\partial C}{\partial s}(s_i) \Delta s + \frac{1}{2} \frac{\partial^2 C}{\partial s^2}(s_i) \Delta s^2 - \frac{1}{3!} \frac{\partial^3 C}{\partial s^3} \Delta s^3 + O(\Delta s^4)$$

$$(1) - (2) \Rightarrow C(s_i + \Delta s) - C(s_i - \Delta s) = 2 \frac{\partial C}{\partial s}(s_i) \Delta s + O(\Delta s^3)$$

the centralized derivative $\frac{\partial C_i^j}{\partial s} = \frac{C_{i+1}^j - C_{i-1}^j}{2 \Delta s} + O(\Delta s^2)$ improves the error

second derivative: (1)+(2)

$$\Rightarrow C(s_i + \Delta s) + C(s_i - \Delta s) = 2C(s_i) + \frac{\partial^2 C}{\partial s^2}(s_i) \Delta s^2 + O(\Delta s^4)$$

$$\Rightarrow \frac{\partial^2 C_i^j}{\partial s^2} = \frac{C_{i+1}^j - 2C_i^j + C_{i-1}^j}{\Delta s^2} + O(\Delta s^2)$$

↳ also $O(\Delta s^2)$ error, as for centralized first derivative

4.4 Stability of Time-stepping Methods

Stability = convergence to true solution

We just consider the simple example of exponential decay

$$\frac{dy}{dt} = -\lambda y, \lambda > 0 \Rightarrow \text{solution: } y(t) = y_0 e^{-\lambda t}$$

we consider also $\lambda \gg 1$

there are two ways to solve this ODE:

- Explicit Euler method: $\frac{y^{j+1} - y^j}{\Delta t} = -\lambda y^j$ (r.h.s. evaluated at j)

$$\Rightarrow y^{j+1} = -\lambda y^j \Delta t + y^j = (1 - \lambda \Delta t) y^j$$

$$\Rightarrow y^M = (1 - \lambda \Delta t)^M y_0 \quad (\text{note: } \lim_{M \rightarrow \infty} y^M = \lim_{M \rightarrow \infty} (1 - \lambda \frac{\Delta t}{M})^M y_0 = e^{-\lambda \Delta t} y_0)$$

We know that $y^M \rightarrow 0$ for large T

This gives us a condition for convergence, i.e., stability: $|1 - \lambda \Delta t| < 1$

↳ need $1 - \lambda \Delta t < 1$, i.e., $\lambda > 0$ ✓

↳ need $-1 + \lambda \Delta t < 1$, i.e., $\Delta t < \frac{2}{\lambda}$

\Rightarrow only for small enough Δt is the discretization stable

- Implicit Euler method: $\frac{y^{j+1} - y^j}{\Delta t} = -\lambda y^{j+1}$ (r.h.s. evaluated at $j+1$)

$$\Rightarrow (1 + \lambda \Delta t) y^{j+1} = y^j \Rightarrow y^{j+1} = \left(\frac{1}{1 + \lambda \Delta t} \right) y^j$$

$$\Rightarrow y^M = \left(\frac{1}{1 + \lambda \Delta t} \right)^M y_0$$

now stability condition is $\left| \frac{1}{1 + \lambda \Delta t} \right| < 1$, which always holds here, since $\lambda > 0$

\Rightarrow implicit scheme is unconditionally stable

4.5 Application to Heat Equation

consider $\frac{\partial V}{\partial t} = \frac{\partial^2 V}{\partial x^2}$, initial value $V(x, 0)$ (want to know $V(x, T)$)
 (for BS: backwards)

take boundary conditions at $x_0=0$ and $x_{\max}=x_{n+1}$

we have: $\frac{\partial V}{\partial t}(x_i, t_j) = \frac{V(x_i, t_{j+1}) - V(x_i, t_j)}{\Delta t} + O(\Delta t)$

$$\frac{\partial^2 V}{\partial x^2}(x_i, t) = \frac{V(x_{i+1}, t) - 2V(x_i, t) + V(x_{i-1}, t)}{\Delta x^2} + O(\Delta x^2)$$

↳ $t = t_j$: explicit scheme

↳ $t = t_{j+1}$: implicit scheme

denote again $V(x_i, t_j) = V_i^j$

explicit: $\frac{V_i^{j+1} - V_i^j}{\Delta t} = \frac{V_{i+1}^j - 2V_i^j + V_{i-1}^j}{\Delta x^2}$

$$\Rightarrow V_i^{j+1} = \frac{\Delta t}{\Delta x^2} V_{i+1}^j + \left(1 - \frac{\Delta t}{\Delta x^2}\right) V_i^j + \frac{\Delta t}{\Delta x^2} V_{i-1}^j$$

We will see numerically that $\frac{\Delta t}{\Delta x^2} < \text{const}$ is needed for stability

implicit: $\frac{V_i^{j+1} - V_i^j}{\Delta t} = \frac{V_{i+1}^{j+1} - 2V_i^{j+1} + V_{i-1}^{j+1}}{\Delta x^2}$

$$\Rightarrow V_i^{j+1} = -\underbrace{\frac{\Delta t}{\Delta x^2}}_{\alpha} V_{i+1}^{j+1} + \left(1 + \frac{\Delta t}{\Delta x^2}\right) V_i^{j+1} - \frac{\Delta t}{\Delta x^2} V_{i-1}^{j+1}$$

In matrix notation:

here boundary eq.s are still wrong

$$\begin{matrix} (l+2a) & -a & & \\ -a & (l+2a) & -a & \\ & \ddots & \ddots & \\ & -a & \ddots & \ddots \\ & & \vdots & \\ & 0 & & \end{matrix} = \begin{pmatrix} v_1^{j+1} \\ v_2^{j+1} \\ \vdots \\ v_n^{j+1} \end{pmatrix} = \begin{pmatrix} v_1^j \\ v_2^j \\ \vdots \\ v_n^j \end{pmatrix}$$

here boundary eq.s are still wrong

\Rightarrow need to solve tridiagonal system of equations to get \vec{V}^{j+1} from \vec{V}^j

what happens at the boundary?

V_o^{j+1} and V_{uti}^{j+1} are given by fixed boundary conditions!

$$\text{we have } V_1^j = -\alpha V_2^{j+1} + (1-\alpha) V_1^{j+1} - \alpha V_0^{j+1}$$

$$V_u^j = -\alpha V_{u+1}^{j+1} + (1+\alpha) V_u^{j+1} - \alpha V_{u-1}^{j+1}$$

so with boundary conditions the tridiagonal system is

$$\left(\begin{array}{cccc} & & & -q \\ & \ddots & & \\ & & (1+2\alpha) & -q \\ -q & & & \ddots \\ & \ddots & & \end{array} \right) \left(\begin{array}{c} V_1^{j+1} \\ \vdots \\ V_n^{j+1} \end{array} \right) = \left(\begin{array}{c} V_1^j + \alpha V_0^{j+1} \\ V_2^j \\ \vdots \\ V_{n-1}^j \\ V_n + \alpha V_{n+1}^{j+1} \end{array} \right)$$

fixed

$$\Rightarrow A \vec{V}^{j+1} = \vec{V}^j + \begin{pmatrix} \alpha V_0^{j+1} \\ 0 \\ \vdots \\ 0 \\ \alpha V_{n+1}^{j+1} \end{pmatrix}$$

tridiagonal

↳ stable scheme with right boundary conditions