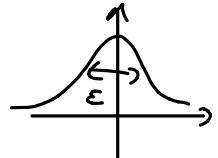


Notes on nonlinear SPDEs:

$$\phi^4 \text{ theory: } \underbrace{\partial_t \phi}_{\text{heat eq.}} = \Delta \phi - \phi^3 + \xi \xleftarrow{\text{white noise}}, \text{ SPDE}$$

\downarrow non-linearity

$$\text{regularization by convolution } \xi_\varepsilon = \xi * \rho_\varepsilon, \quad \int \rho_\varepsilon = \int \rho = 1$$



$$\Rightarrow \partial_t \phi_\varepsilon = \Delta \phi_\varepsilon - \phi_\varepsilon^3 + \xi_\varepsilon \quad \text{can be solved as classical PDE} \\ (\phi_\varepsilon \text{ differentiable})$$

but $\phi_\varepsilon \rightarrow 0$ as $\varepsilon \rightarrow 0$ (can be shown in dim. d=2,3)

introducing a counterterm $C_\varepsilon \phi_\varepsilon$ in equation can make new $\phi_\varepsilon \xrightarrow{\varepsilon \rightarrow 0} \phi$

\hookrightarrow idea of renormalization of an SPDE

vigorous solution theory: Martin Hairer (Fields medal 2014)

Notes on the connection from SDE to PDE:

this goes via the distribution fct.:

1to SDE: $dX = f(x)dt + g(x)dW$

$$\begin{aligned} \mathbb{E}(F(X_t)) &= \int_{-\infty}^{\infty} \rho(x, t) F(x) dx \\ \implies \frac{d\mathbb{E}(F(x_t))}{dt} &= \int_{-\infty}^{\infty} \frac{\partial \rho(x, t)}{\partial t} F(x) dx \end{aligned}$$

$$\begin{aligned} \text{with 1to's lemma: } \mathbb{E}(dF(X_t)) &= \mathbb{E}\left(\frac{dF}{dx}(f(x)dt + g(x)dW) + \frac{1}{2} \frac{d^2F}{dx^2} g(x)^2 dt\right) \\ &= \int_{-\infty}^{\infty} \rho(x, t) \left(\frac{dF}{dx} f(x) + \frac{1}{2} \frac{d^2F}{dx^2} g(x)^2 \right) dx dt \\ &\stackrel{\text{integration by parts}}{=} \int_{-\infty}^{\infty} \left[-\frac{\partial}{\partial x} (\rho(x, t) f(x)) + \frac{\partial^2}{\partial x^2} (\rho(x, t) \frac{1}{2} g(x)^2) \right] F(x) dx dt \end{aligned}$$

Comparing the expressions yields

$$\frac{\partial \rho}{\partial t} = \frac{1}{2} \frac{\partial^2}{\partial x^2} (g^2 \rho) - \frac{\partial}{\partial x} (f \rho) \quad , \text{ Fokker-Planck equation}$$

Ex.: Ornstein-Uhlenbeck process: $dX = -\alpha V'(x)dt + \sqrt{2\beta} dW$

$$\implies \frac{\partial \rho}{\partial t} = \beta \frac{\partial^2 \rho}{\partial x^2} + \alpha \frac{\partial}{\partial x} (V' \rho)$$

if there is equilibrium for large t , then $\frac{\partial \rho}{\partial t} = 0$ (large t)

$$\Rightarrow \frac{\partial}{\partial x} \left(\beta \frac{\partial \rho_e}{\partial x} + \alpha V' \rho_e \right) = 0$$

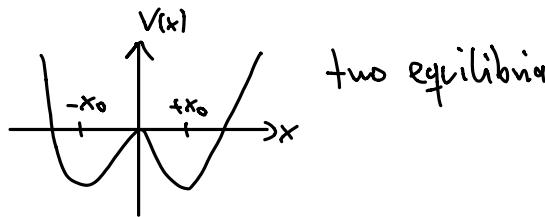
Suppose $\frac{\partial \rho_e}{\partial x}, V', \rho_e \xrightarrow{x \rightarrow \pm\infty} 0$

$$\Rightarrow \text{need to solve } \beta \frac{\partial \rho_e}{\partial x} + \alpha V' \rho_e = 0$$

$$\Rightarrow \frac{\partial \rho_e}{\partial x} = -\frac{\alpha}{\beta} V' \rho_e \quad \text{equilibrium}$$

$$\Rightarrow \frac{d \rho_e}{\rho_e} = -\frac{\alpha}{\beta} V' dx \Rightarrow \rho_e(x) = C e^{-\frac{\alpha}{\beta} V(x)}$$

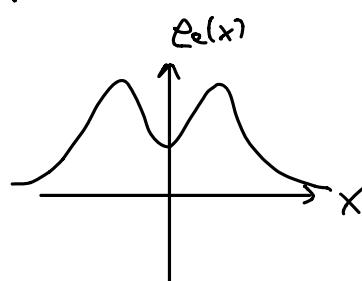
e.g., for $V(x) = \frac{1}{4}x^4 - \frac{1}{2}x^2$:



without noise ($\eta=0$), the solution would end up either in $-x_0$ or in $+x_0$.

with noise: switching between equilibria

(distribution fct.)



e.g., useful for detecting weak signals in noise