

3.3 Line Integrals

Next: We consider integrals along curves and surfaces and their relation to Riemann integrals and each other. This will lead us to generalizations of the Fundamental Theorem of Calculus. (E.g., $\int_{\text{curve } \gamma} F \, dx$ depends only on endpoints $x(a)$ and $x(b)$. E.g., $\int_D \nabla \cdot b \, dS = \int_D b \cdot n \, ds.$)

integral over D depends only on ∂D !

Applications: Force fields, electrodynamics, ...

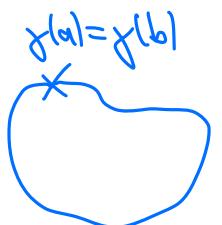
Definition:

Any continuous function $\gamma: [a, b] \rightarrow \mathbb{R}^n$ is called an **oriented curve** (or path).

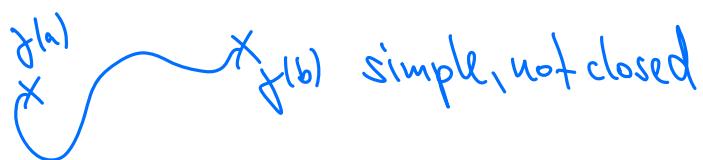


First, a few important types of curves:

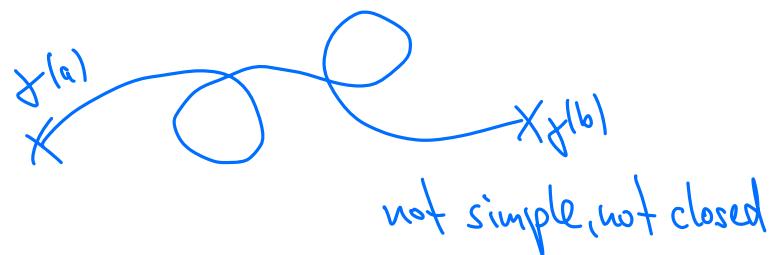
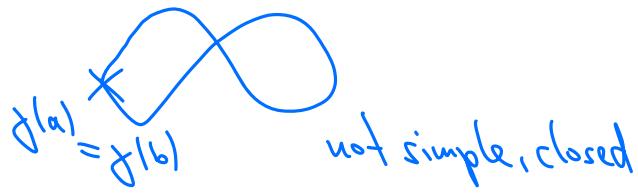
- If $\gamma(a) = \gamma(b)$, the curve is **closed**.
- If $\gamma: [a, b] \rightarrow \mathbb{R}^n$ is injective, the curve is **simple**.



simple closed curve



simple, not closed



- Two curves f and g are called **equivalent** if there is a continuous, monotonic, increasing h s.t. $f = g \circ h$ (i.e., the images of f and $g \circ h$ are the same).
- "going through the curve with a different velocity"

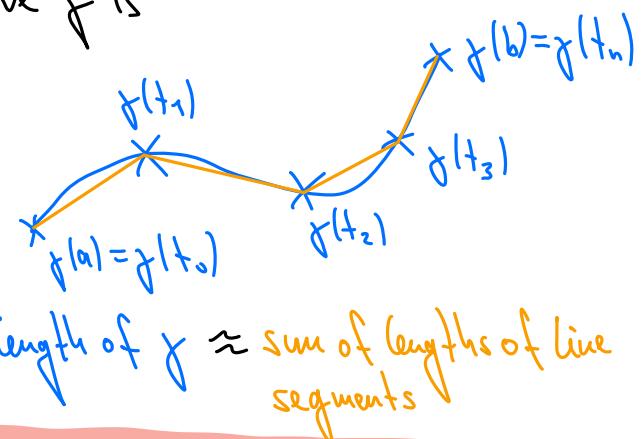
Next, we define the length of a curve (not all curves have a length!)

Let \mathcal{T} be a partition of $[a, b]$ with $a = t_0 < t_1 < t_2 < \dots < t_n = b$, and let

$$\lambda(\mathcal{T}) := \max_{i=1, \dots, n} \underbrace{|t_i - t_{i-1}|}_{=: \Delta t_i}$$

Then an approximation to the length of a curve f is

$$L(\mathcal{T}, f) := \sum_{i=1}^n \|f(t_i) - f(t_{i-1})\|.$$



Definition: The **length of the curve** $f \in C([a, b])$ is defined as

$$L(f) = \sup_{\mathcal{T}} L(\mathcal{T}, f).$$

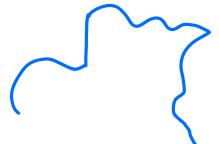
If $L(f) < \infty$, we call f **rectifiable** (" f has length").

We get a more concrete formula for continuously differentiable curves.

Theorem: Let $\gamma \in C^1([a,b])$. Then γ is rectifiable and

$$L(\gamma) = \int_a^b \|\gamma'(t)\| dt.$$

Note: The theorem is obviously extended to piece-wise C^1 curves.



Proof:

$$\begin{aligned} " \leq ": L(\gamma, \tau) &= \sum_{i=1}^n \|\gamma(t_i) - \gamma(t_{i-1})\| = \sum_{i=1}^n \left\| \int_{t_{i-1}}^{t_i} \gamma'(t) dt \right\| \\ &\leq \sum_{i=1}^n \int_{t_{i-1}}^{t_i} \|\gamma'(t)\| dt = \int_a^b \|\gamma'(t)\| dt, \\ \text{so also } L(\gamma) &= \sup_{\tau} L(\gamma, \tau) \leq \int_a^b \|\gamma'(t)\| dt. \end{aligned}$$

" \geq ": Let $\epsilon > 0$. We know that γ' is uniformly continuous $\Rightarrow \exists \delta > 0$ s.t. $\forall s, t \in [a, b]$ with $|s-t| < \delta$ we have $\|\gamma(s) - \gamma(t)\| < \epsilon$.

Let τ be a partition with $\lambda(\tau) < \delta$.

$$\text{Then } \|\gamma'(t)\| \leq \|\gamma'(t_i)\| + \epsilon \quad \forall t \in [t_{i-1}, t_i]$$

$$\Rightarrow \int_a^b \|\gamma'(t)\| dt = \sum_{i=1}^n \int_{t_{i-1}}^{t_i} \|\gamma'(t)\| dt \leq \sum_{i=1}^n (\|\gamma'(t_i)\| + \epsilon) \Delta t_i$$

$$= \sum_{i=1}^n \left(\underbrace{\left\| \int_{t_{i-1}}^{t_i} \gamma'(t) dt \right\|}_{\text{length of subinterval}} + \epsilon \Delta t_i \right)$$

$$= \int_a^b \gamma'(t) dt + \sum_{i=1}^n (\gamma'(t_i) - \gamma'(t_{i-1})) \Delta t_i \xrightarrow{\leq \epsilon \text{ in abs. value}}$$

$$\begin{aligned}
 &\leq \sum_{i=1}^n \left(\underbrace{\left\| \int_{t_{i-1}}^{t_i} \gamma'(t) dt \right\|}_{\gamma(t_i) - \gamma(t_{i-1})} + 2\varepsilon \Delta t_i \right) \\
 &= \underbrace{\Lambda(\gamma)}_{\gamma(t_i) - \gamma(t_{i-1})} + 2\varepsilon(b-a) \\
 &\leq \Lambda(\gamma)
 \end{aligned}$$

Since ε was arbitrary (arbitrarily small), we find $\int_a^b \| \gamma'(t) \| dt \leq \Lambda(\gamma)$. \square

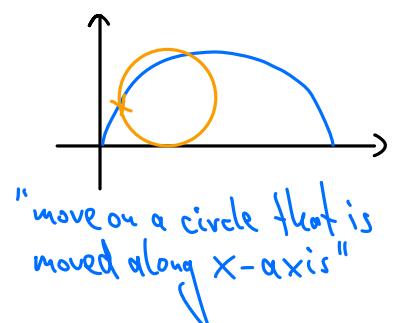
Note that for $\gamma \in C^1$, the length $\Lambda(\gamma)$ is independent of the parameterization: If $\gamma = \rho \circ h$, $h \in C^1$ increasing, then

$$\int_a^b \| \gamma'(t) \| dt = \int_a^b \left\| \frac{d}{dt} \rho(h(t)) \right\| dt = \int_a^b \| \rho'(h(t)) \| |h'(t)| dt = \int_{h(a)}^{h(b)} \| \rho'(u) \| du.$$

chain rule substitution
($u = h(t) \Rightarrow du = h'(t) dt$)

Examples:

- Non-rectifiable curve: see homework.
- Circumference of a circle: $\gamma(t) = R(\cos t, \sin t)$, $t \in [0, 2\pi]$
 $\Rightarrow \gamma'(t) = R(-\sin t, \cos t) \Rightarrow \| \gamma'(t) \| = \sqrt{R^2 \sin^2 t + R^2 \cos^2 t} = R$.
 $\Rightarrow \Lambda(\gamma) = \int_0^{2\pi} R dt = 2\pi R$.
- Cycloid: $\gamma(t) = (t - \sin t, 1 - \cos t)$, $t \in [0, 2\pi]$
 $\Rightarrow \gamma'(t) = (1 - \cos t, \sin t)$



$$\Rightarrow \|\gamma'(t)\|^2 = (1-\cos t)^2 + \sin^2 t = 1 - 2\cos t + \cos^2 t + \sin^2 t = 2(1-\cos t)$$

standard trigonometric identity

$$= 4 \sin^2\left(\frac{t}{2}\right)$$

$$\Rightarrow A(\gamma) = \int_0^{2\pi} 2 \sin\left(\frac{t}{2}\right) dt = -4 \cos\left(\frac{t}{2}\right) \Big|_0^{2\pi} = 4 - (-4) = 8.$$

One useful parametrization is the arc length parametrization:

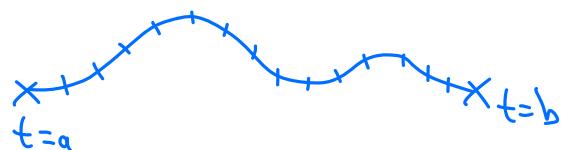
Define $s(t) = \int_a^t \|\gamma'(\tau)\| d\tau$. Then $s'(t) = \|\gamma'(t)\| > 0$ for $\gamma'(t) \neq 0$
 ("non-degenerate parametrization").

Since $s(t)$ is monotonic it is invertible, with inverse $t(s)$.

We call $\gamma(t(s))$ the arc length parametrization.

$$\Rightarrow \frac{d}{ds} \gamma(t(s)) = \gamma'(t(s)) \frac{dt(s)}{ds} = \gamma'(t(s)) \frac{1}{s'(t(s))} \Rightarrow \left\| \frac{d}{ds} \gamma(t(s)) \right\| = \left\| \frac{\gamma'}{s'} \right\| = 1,$$

i.e., we go through the curve with speed 1.



Definition: For $f \in C(\gamma, \mathbb{R})$, and γ a C^1 curve, we define the

line integral

$$\int_{\gamma} f ds = \int_0^{A(\gamma)} f(\gamma(t(s))) ds.$$

Note: $\int_{\gamma} f ds := \int_0^{A(\gamma)} f(\gamma(t(s))) ds = \int_a^b f(\gamma(t)) \|\gamma'(t)\| dt.$

Substitution $t = t(s) \Rightarrow \frac{dt}{ds} = \frac{1}{s'(t)} = \frac{1}{\|\gamma'(t)\|} \Rightarrow ds = \|\gamma'(t)\| dt$

Note also that $\int_{\gamma} f ds$ is independent of the parametrization of γ (see HW).

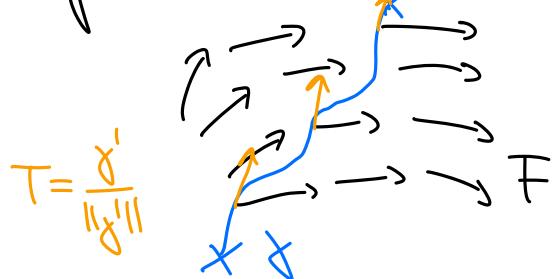
Next: How to def. line integrals for vector fields $F: D \subset \mathbb{R}^n \rightarrow \mathbb{R}^n$?
 (E.g., force fields in physics.)

In physics: work = force \cdot length.
 or rather "displacement"

In 1-dim.: $W = \int F(s) ds$.

In n -dimensions: Only displacement in the direction of the force is work
 (e.g., displacement orthogonal to force causes no work).

In general: If T is the unit tangent vector, then work = $\langle F, T \rangle$.



$$\begin{aligned} \Rightarrow \text{Total work } W &= \int_{\gamma} \langle F, T \rangle ds = \int_{\gamma} \langle F, \frac{\gamma'}{\|\gamma'\|} \rangle ds \\ &= \int_a^b \langle F(\gamma(t)), \frac{\gamma'(t)}{\|\gamma'(t)\|} \rangle \|\gamma'(t)\| dt \\ &= \int_a^b \langle F \circ \gamma, \gamma' \rangle dt \\ &= \int_a^b F(\gamma(t)) \cdot \gamma'(t) dt. \end{aligned}$$

Definition: For $F \in C(\gamma, \mathbb{R}^n)$ and γ a C^1 curve, we define

the line integral $\int_{\gamma} F ds = \int_a^b F(\gamma(t)) \gamma'(t) dt$.

Note: $F dx := F_1 dx_1 + \dots + F_n dx_n$ is called a first-order differential form.