

An example of an infinite dim. symplectic vector space.

Let \mathcal{H} be a Hilbert space, let us take $\mathcal{H} = L^2(\mathbb{T}^d)$ to be specific.

\Rightarrow scalar product $\langle f, g \rangle = \int \overline{f(x)} g(x) dx$.

Now let \bar{f} be def. by $\bar{f}(x) = \overline{f(x)}$ (complex conjugation), and def. the vector space $\mathcal{Y} = \{(f, \bar{f}), f \in \mathcal{H}\}$.

We def. $\epsilon((f, \bar{f}), (g, \bar{g})) = \text{Im} \langle f, g \rangle$

\hookrightarrow nondegenerate

$\hookrightarrow \epsilon((g, \bar{g}), (f, \bar{f})) = \text{Im} \langle g, f \rangle = \text{Im} \overline{\langle f, g \rangle} = -\text{Im} \langle f, g \rangle = -\epsilon((f, \bar{f}), (g, \bar{g}))$

$\Rightarrow (\mathcal{Y}, \epsilon)$ is a symplectic vector space

Note:

$$\langle (f, \bar{f}), \frac{1}{2i} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} g \\ \bar{g} \end{pmatrix} \rangle = \frac{1}{2i} (\langle f, g \rangle - \langle \bar{f}, \bar{g} \rangle) = \frac{1}{2i} (\langle f, g \rangle - \langle g, f \rangle) = \text{Im} \langle f, g \rangle = \epsilon((f, \bar{f}), (g, \bar{g})).$$

If $R: \mathcal{Y} \rightarrow \mathcal{Y}$ is symplectic, then $\epsilon(R(f, \bar{f}), R(g, \bar{g})) = \epsilon((f, \bar{f}), (g, \bar{g}))$.

$$\Rightarrow R^* \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} R = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Let us consider specifically $R = \begin{pmatrix} u & \bar{v} \\ v & \bar{u} \end{pmatrix}$, with $u, v: \mathcal{H} \rightarrow \mathcal{H}$ (bounded), $\bar{u}f := \overline{uf}$.

$$\text{Then } R^* \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} R = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \Leftrightarrow \underbrace{\begin{pmatrix} u^* & \bar{v}^* \\ \bar{v}^* & \bar{u}^* \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} u & \bar{v} \\ v & \bar{u} \end{pmatrix}}_{= \begin{pmatrix} u & \bar{v} \\ -v & -\bar{u} \end{pmatrix}} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\Rightarrow R \text{ is symplectic} \iff u^*u - v^*v = 1, \bar{u}^*\bar{u} - \bar{v}^*\bar{v} = 1, \quad \left. \begin{array}{l} u^*\bar{v} - v^*\bar{u} = 0, \\ \bar{v}^*u - \bar{u}^*v = 0 \end{array} \right\} (*)$$

An R that satisfies $(*)$ is called Bogoliubov map.

A few properties of Bogoliubov maps:

- $u^*u = 1 + v^*v \geq 1 \Rightarrow u$ is invertible

- Specific example for $L^2([0,1])$: $u = \sum_{k=-\infty}^{\infty} g_k |e^{ikx}\rangle \langle e^{ikx}|, g_k \in \mathbb{R}$

$$v = \sum_{k=-\infty}^{\infty} (-j_k) |e^{ikx}\rangle \langle e^{ikx}|, j_k \in \mathbb{R}$$

with $|g_k|^2 - |j_k|^2 = 1, |g_k| \geq 1$, i.e., $g_k = \cosh \lambda_k, j_k = \sinh \lambda_k$ for some $\lambda_k \in \mathbb{R}$.

Recall: A linear map $f: V \rightarrow V$ is symplectic if and only if $A^T J A = J$, where A is the matrix of f and J is the matrix of the symplectic form w of V .

Now: Suppose f (with matrix A) and g (with matrix B) are symplectic.

Then $(AB)^T J (AB) = B^T \underbrace{A^T J A}_= B^T J B = J$, so also AB is symplectic.

\Rightarrow Symplectic maps form a group

Recall:

- Def.: • A 2-form w on a smooth manifold M is called nondegenerate if w_p is a nondegenerate 2-covector for each $p \in M$. exterior derivative = 0
- A symplectic form on M is a closed nondegenerate 2-form.
- M with a symplectic form is called symplectic manifold.
- or symplectic structure

Note: • From the discussion above we know that $\dim M$ must be even.

- Also: If w is a symplectic form on a $2n$ -dim. manifold M , then w^n is a nonvanishing $2n$ -form, so every symplectic manifold is orientable.
- If (M_1, w_1) and (M_2, w_2) are symplectic manifolds, then a diffeomorphism $F: M_1 \rightarrow M_2$ with $F^*w_2 = w_1$ is called symplectomorphism \Rightarrow symplectic geometry/topology

Ex.: • \mathbb{R}^{2n} with coordinates $(x^1, \dots, x^n, y^1, \dots, y^n)$:

$$w = \sum_{i=1}^n dx^i \wedge dy^i \text{ is symplectic } (\text{closed: } \checkmark \text{ nondegenerate: } \checkmark)$$

- Σ an oriented smooth 2-manifold, w nonvanishing 2-form.
 $\Rightarrow w$ closed ($dw = 3\text{-form}$), $w = \alpha^1 \wedge \beta^1$ locally i.e. nondegenerate

Globally, we have the following very strong result.

Theorem (Darboux): Let (M, w) be a $2n$ -dim. symplectic manifold. Then for any $p \in M$ there are smooth coordinates $(x^1, \dots, x^n, y^1, \dots, y^n)$ centered at p in which w has coordinate representation
 $w = \sum_{i=1}^n dx^i \wedge dy^i$.

Proof: Uses time-dependent flows, we skip it here.

Note: Such coordinates are called Darboux/symplectic/canonical coordinates.

Next:

Def.: Given a symplectic manifold (M, ω) and a smooth function $f \in C^\infty(M)$, we def.
the Hamiltonian vector field of f to be X_f , def. via

$$X_f \lrcorner \omega = df, \text{ or, } \omega(X_f, Y) = df(Y) = Yf \text{ for any vector field } Y.$$

In Darboux coordinates:

$$X_f = \sum_{i=1}^n \left(a^i \frac{\partial}{\partial x^i} + b^i \frac{\partial}{\partial y^i} \right), a^i, b^i \text{ some coefficient functions}$$

$$\hookrightarrow X_f \lrcorner \omega = \sum_{i=1}^n \left(a^i \frac{\partial}{\partial x^i} + b^i \frac{\partial}{\partial y^i} \right) \lrcorner \sum_{j=1}^n dx^j \wedge dy^j \\ = \sum_{i=1}^n (a^i dy^i - b^i dx^i)$$

$$\hookrightarrow df = \sum_{i=1}^n \left(\frac{\partial f}{\partial x^i} dx^i + \frac{\partial f}{\partial y^i} dy^i \right)$$

$$\implies a^i = \frac{\partial f}{\partial y^i}, b^i = -\frac{\partial f}{\partial x^i}$$

$$\implies X_f = \sum_{i=1}^n \left(\frac{\partial f}{\partial y^i} \frac{\partial}{\partial x^i} - \frac{\partial f}{\partial x^i} \frac{\partial}{\partial y^i} \right)$$

Note: $X_f f = df(X_f) = \omega(X_f, X_f) = 0$, so f does not change in the direction X_f
 $\implies f$ is constant along integral curves of X_f .

Def.: A symplectic manifold (M, ω) together with $H \in C^\infty(M)$ is called

Hamiltonian system, and

- H is called Hamiltonian
- the flow of X_H is called Hamiltonian flow
- the integral curves of X_H are called trajectories or orbits

Trajectories $\gamma(t) = (x^i(t), \dot{x}^i(t))$ satisfy $\begin{cases} \dot{x}^i(t) = \frac{\partial H}{\partial p_i}(x(t), \dot{x}(t)) \\ \dot{p}_i(t) = -\frac{\partial H}{\partial x^i}(x(t), \dot{x}(t)) \end{cases}$ } Hamilton's eq.s

derivative
with respect
to t

Example: Classical mechanics, n -body problem

- n particles in \mathbb{R}^3 : $(q_1(t), \dots, q_n(t)) \in \mathbb{R}^{3n}$, $q_i(t) \in \mathbb{R}^3$
- configuration space $Q = \mathbb{R}^{3n} \setminus C$, $C = \{q \in \mathbb{R}^{3n}, q_k = q_\ell \text{ for some } k \neq \ell\}$
= collision set
- Newton's eq.s: $m \ddot{q}_k(t) = F_k(q(t))$ for some smooth forces $F \in \mathbb{R}^{3n}$.
- Assume F conservative, take F_k as components of covector field on Q : $F = -dV$,
for some $V \in C^\infty(Q)$
- Consider cotangent space T^*Q with coordinates (q^i, p_i) , then $q(t)$ satisfies
Newton's law if $\dot{\gamma}(t) = (q(t), \dot{p}(t))$ in T^*Q satisfies

$$\dot{q}^i(t) = \frac{1}{m} \sum_{j=1}^n \delta^{ij} p_j(t)$$

$$\dot{p}_i(t) = -\frac{\partial V}{\partial q^i}(q(t))$$

$$\Rightarrow \text{Hamiltonian } H(q, p) = \frac{1}{2m} \dot{p}^2 + V(q)$$

Note: $H = \text{const}$ along flow \Rightarrow conservation of energy.

- Further constraints may lead to considering different configuration spaces Q .